

herramientas de

HPC para

física de colisiones y

reacciones químicas

herramientas de

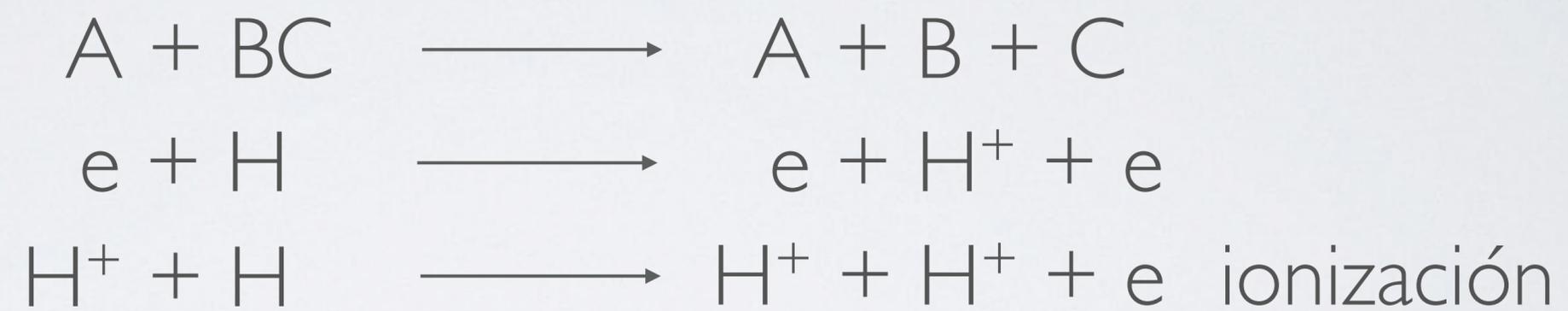
HPC para

física de colisiones y

reacciones químicas

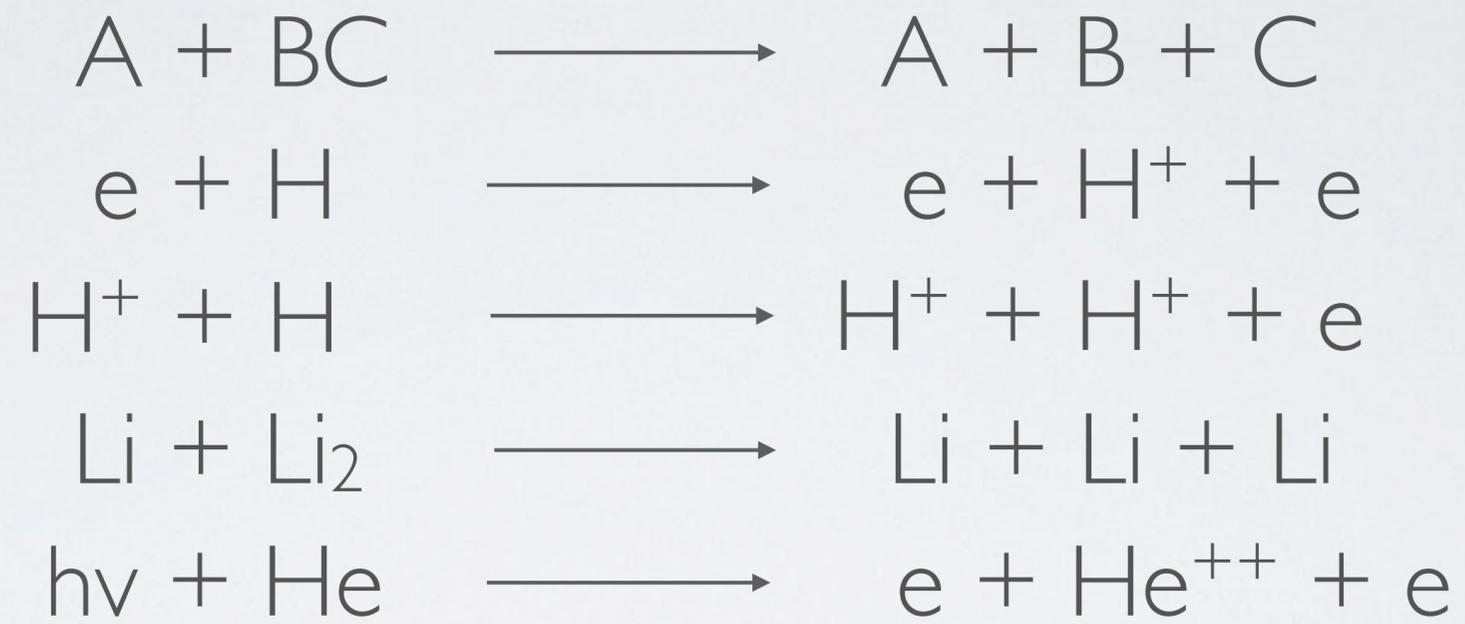








disociación  
inducida por  
colisión



Doble  
Fotoionización



problema de 3 cuerpos cuántico

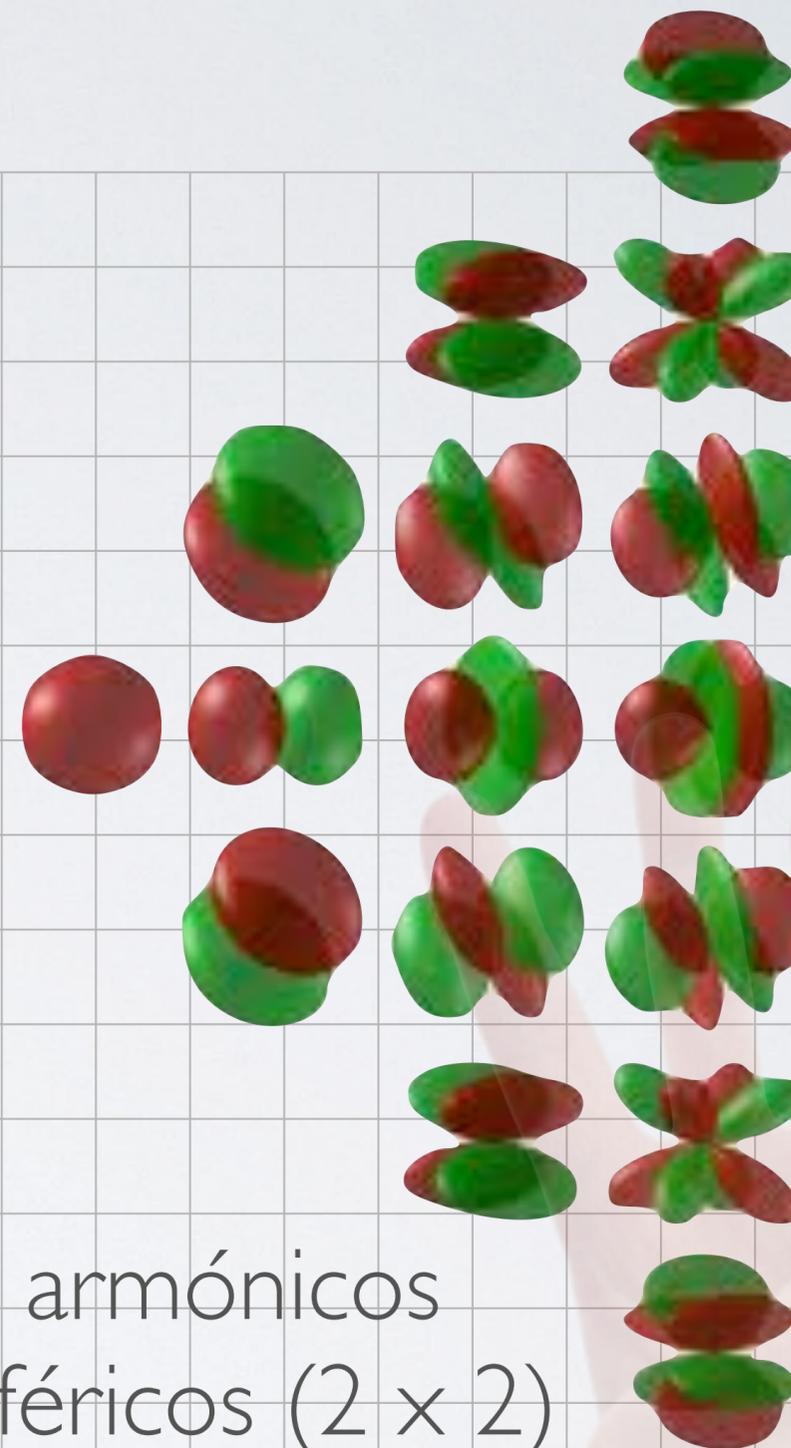
ecuación de  
schrödinger  
independiente del tiempo

$$3 \text{ cuerpos} \times 3 \text{ coord. c/u} - 3 \text{ coord. CM} = 6 \text{ coord}$$

esféricas



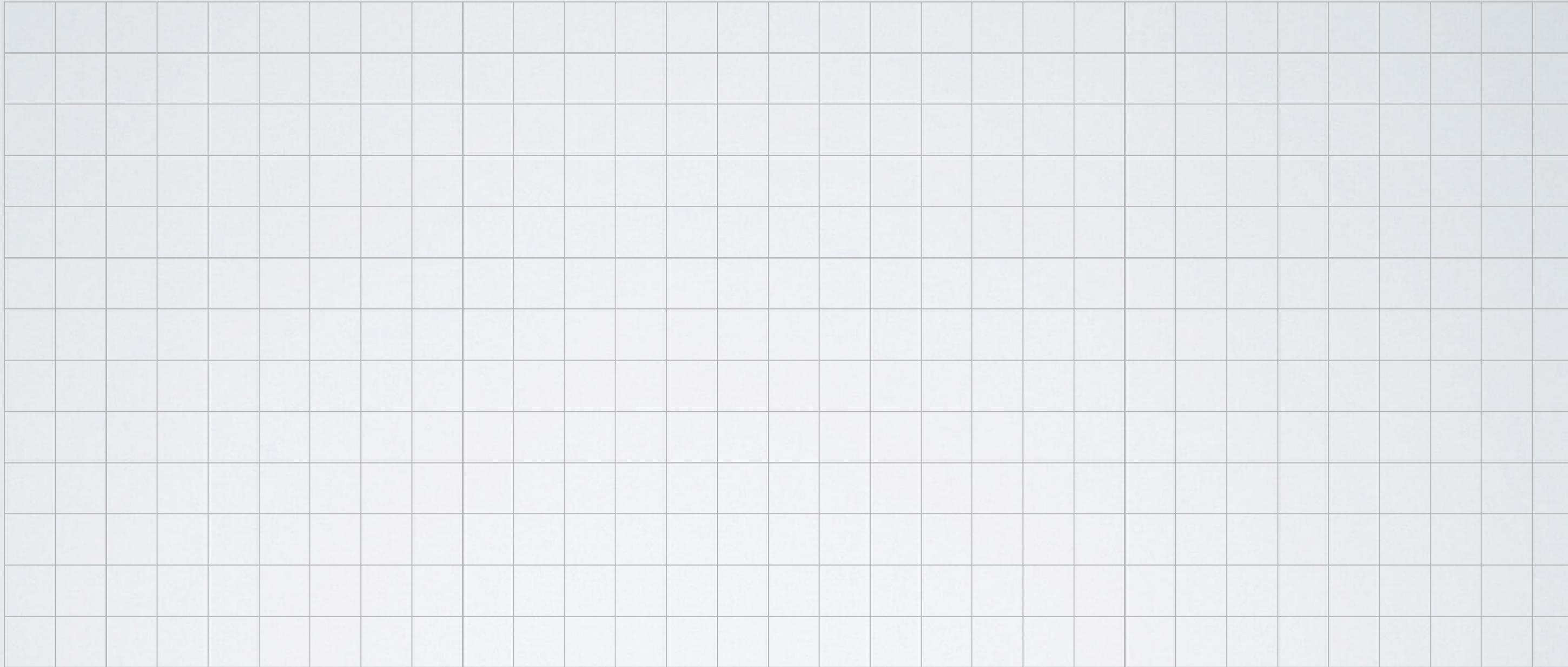
coordenadas esféricas ( $2 \times 3$ )



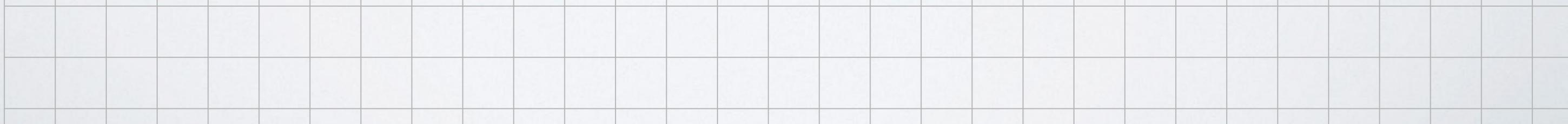
armónicos  
esféricos ( $2 \times 2$ )

dominio radial infinito (2D)

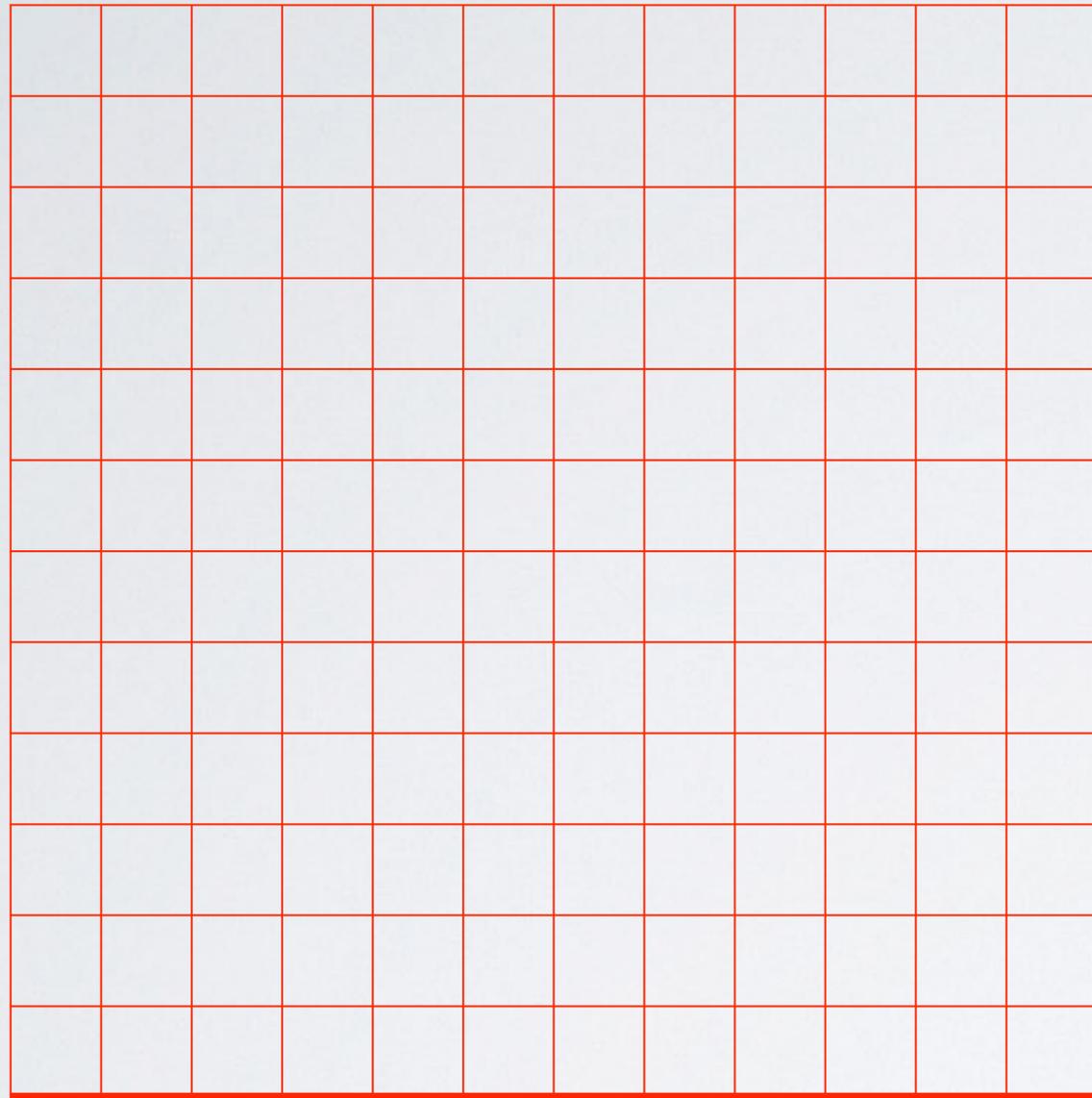
coordenadas esféricas radiales (x2)



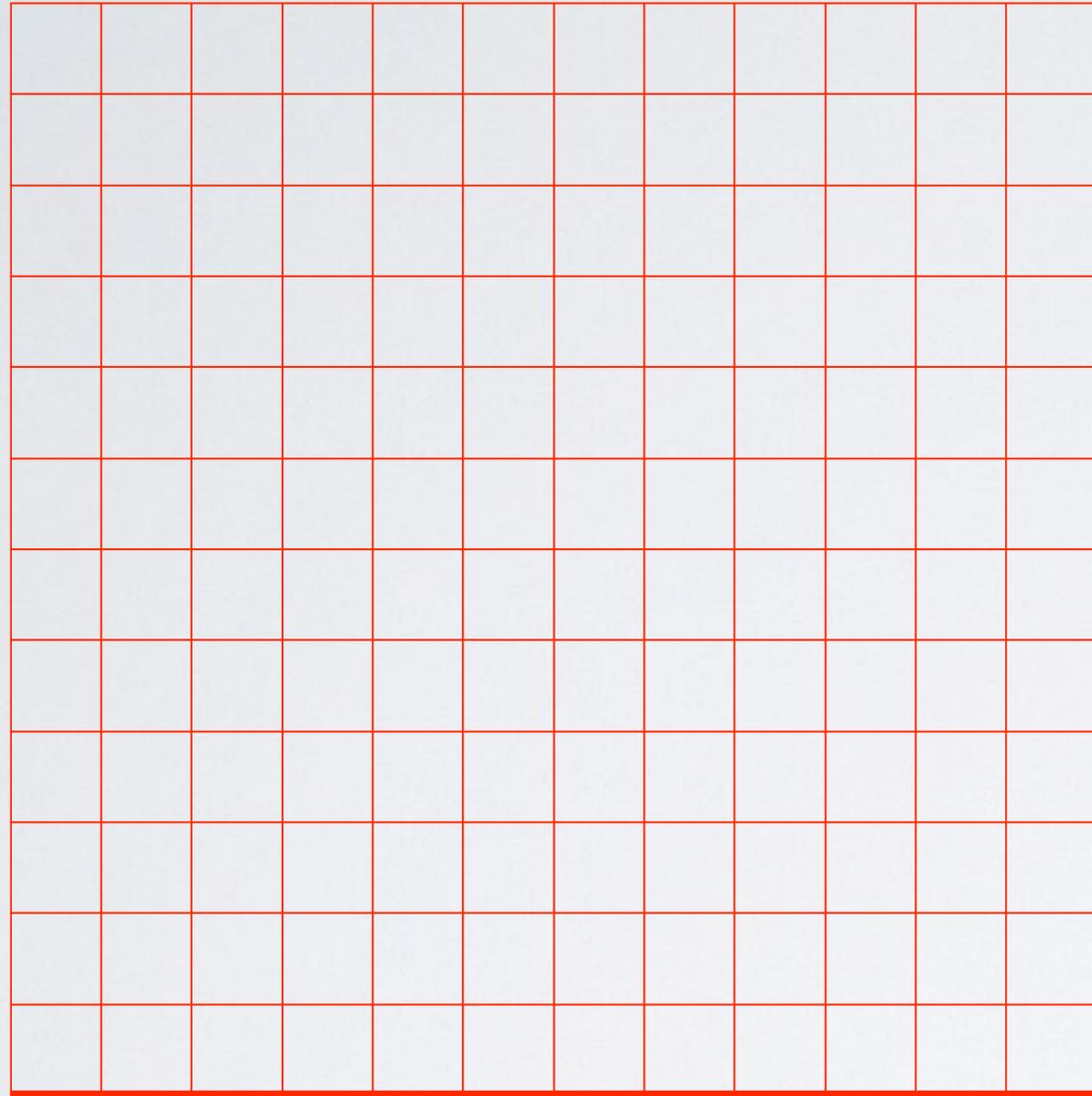
dominio radial infinito (2D)



# funciones sturmianas generalizadas

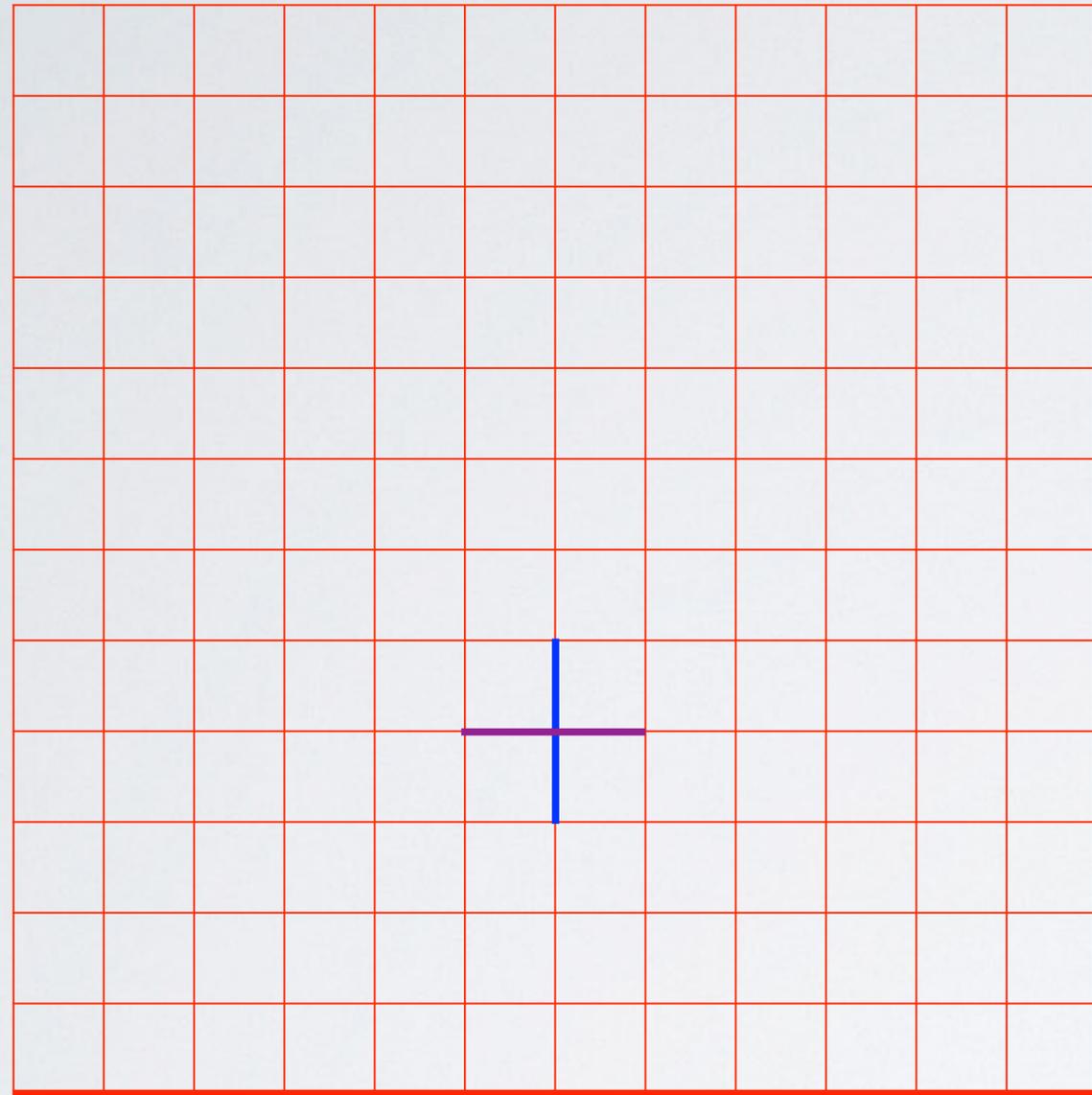


dominio radial **finito** (2D)+ C. C.



ecuación  
diferencial de  
segundo orden en  
derivadas parciales con  
condiciones de  
contorno complejas

dominio radial finito (2D) + C. C.



dominio radial finito (2D) + C. C.

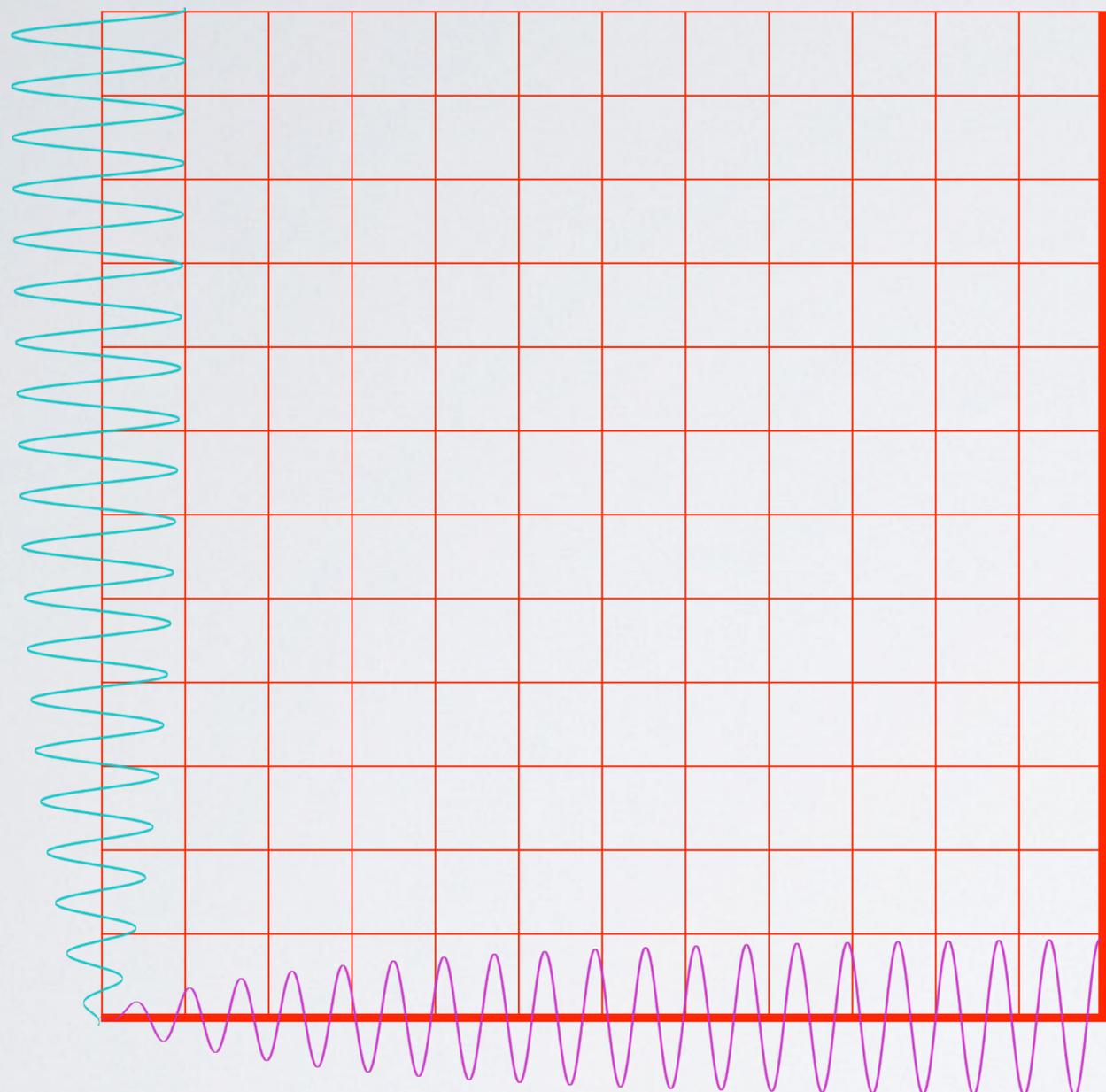
discretización en  
diferencias finitas

cálculo de la  
matriz  $H$



problema de álgebra  
lineal numérica  
ralo





dominio radial finito (2D) + C. C.

método espectral  
con funciones sturmianas

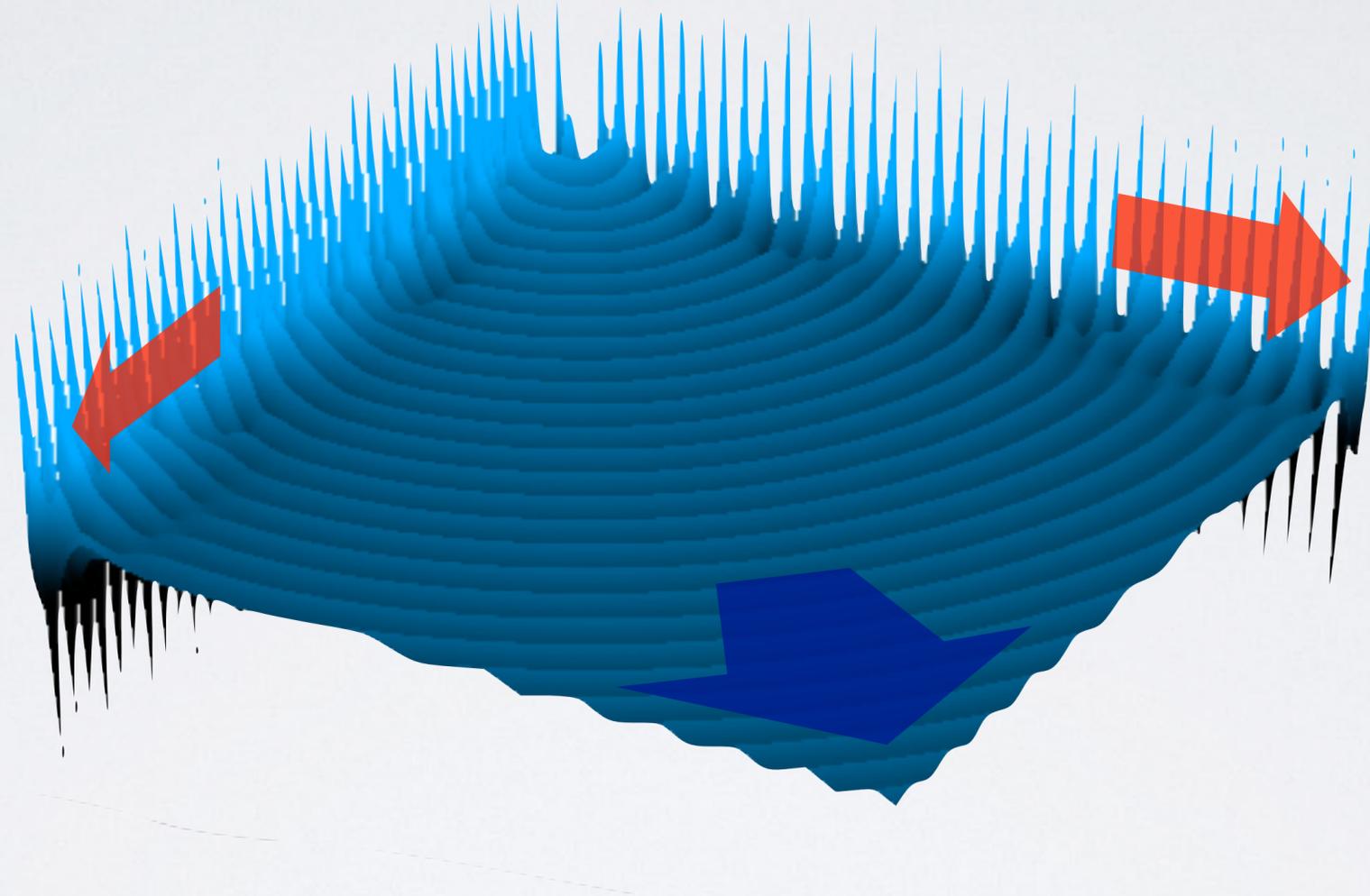
cálculo de la  
matriz  $H$



problema de álgebra  
lineal numérica  
denso



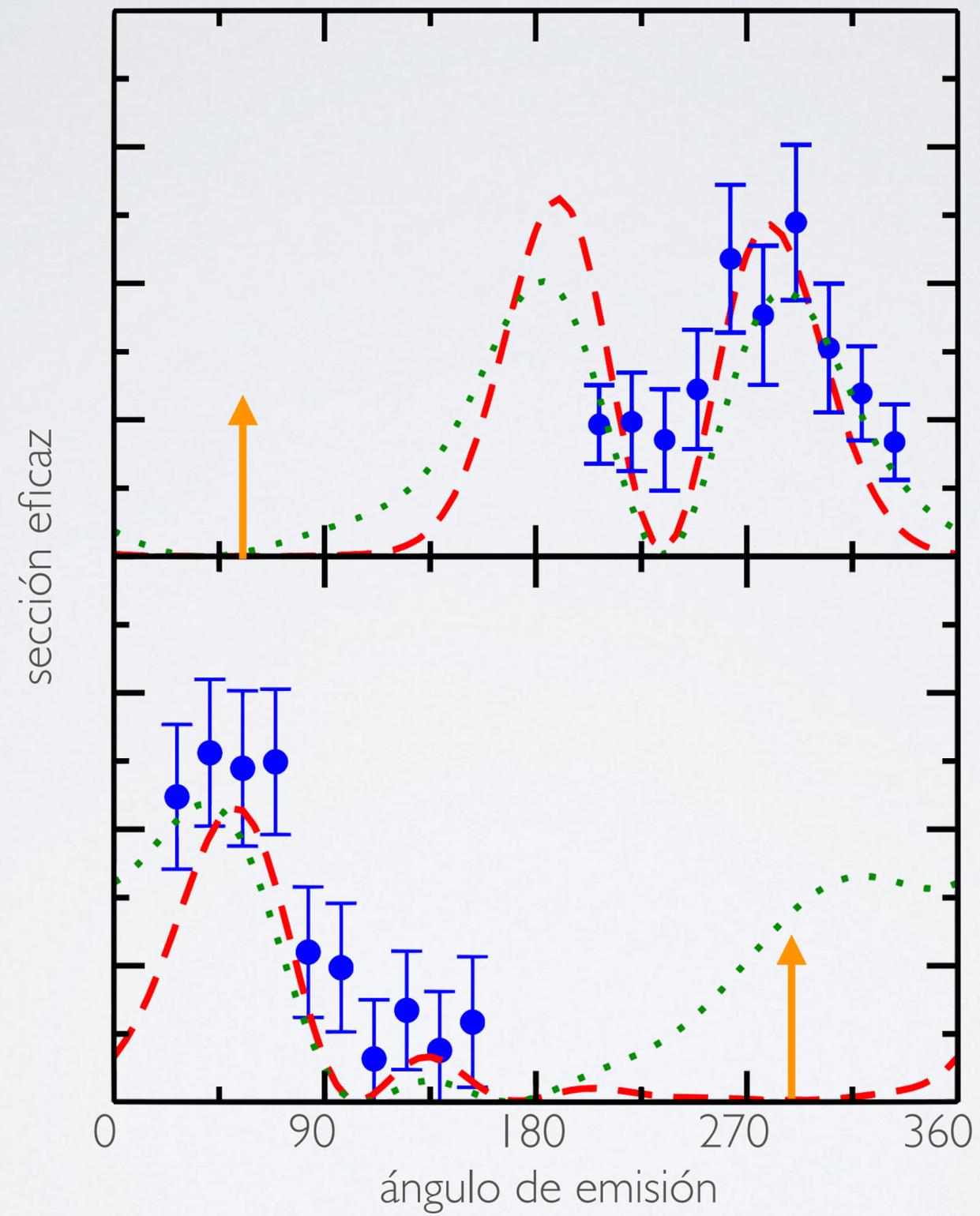
(e,3e) alta energía



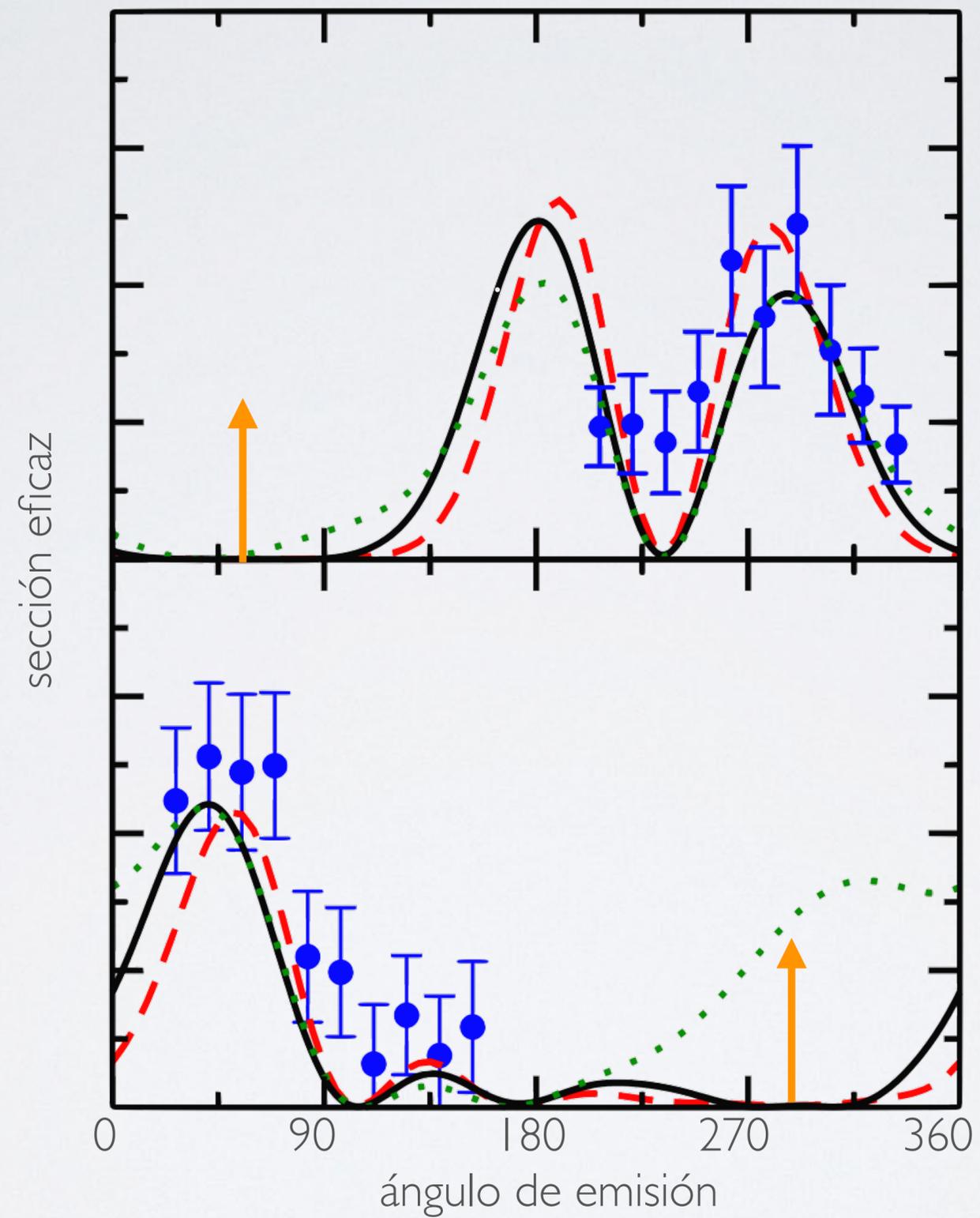
ionización simple  
ionización - excitación

doble ionización

(e,3e) alta energía



(e,3e) alta energía



## Two-body Coulomb wavefunctions as kernel for alternative integral transformations

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Received 16 April 2003, in final form 19 June 2003  
Published 23 July 2003  
Online at [stacks.iop.org/JPhysA/36/8443](http://stacks.iop.org/JPhysA/36/8443)

### Abstract

In this paper we investigate different representations of an arbitrary function in terms of two-body Coulomb eigenfunctions. We discuss the standard energy basis in spherical and parabolic coordinates with the purpose of remarking explicitly that two additional parameters appear both in the Schrödinger equation and in the wavefunctions: the charge and the angular momentum. We introduce the charge and generalized angular momentum Sturmian function representations, which result when the charge or the angular momentum is used as the eigenvalue in the Coulomb Schrödinger equation, respectively. We present the connection between the generalized angular momentum representation and the Kontorovich–Lebedev transform. Finally, we extend the angular momentum representation to six dimensions, which is suitable for further applications in the three-body Coulomb problem.

PACS number: 03.65.–w

## 1. Introduction

The standard way of theoretically studying the properties of a wide variety of atomic and molecular systems is by using functional basis set for representing the physical magnitudes involved in the phenomena under analysis. According to the mathematical structure of quantum mechanics, each physical magnitude has a Hermitian operator associated and then its eigenfunctions become the natural basis set to be used. The most common basis set uses the energy eigenfunctions, i.e., the eigenfunctions of the Hamiltonian. However, in some cases the energy eigenfunctions are not the most convenient ones and alternative basis sets are necessary.

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## FAST TRACK COMMUNICATION

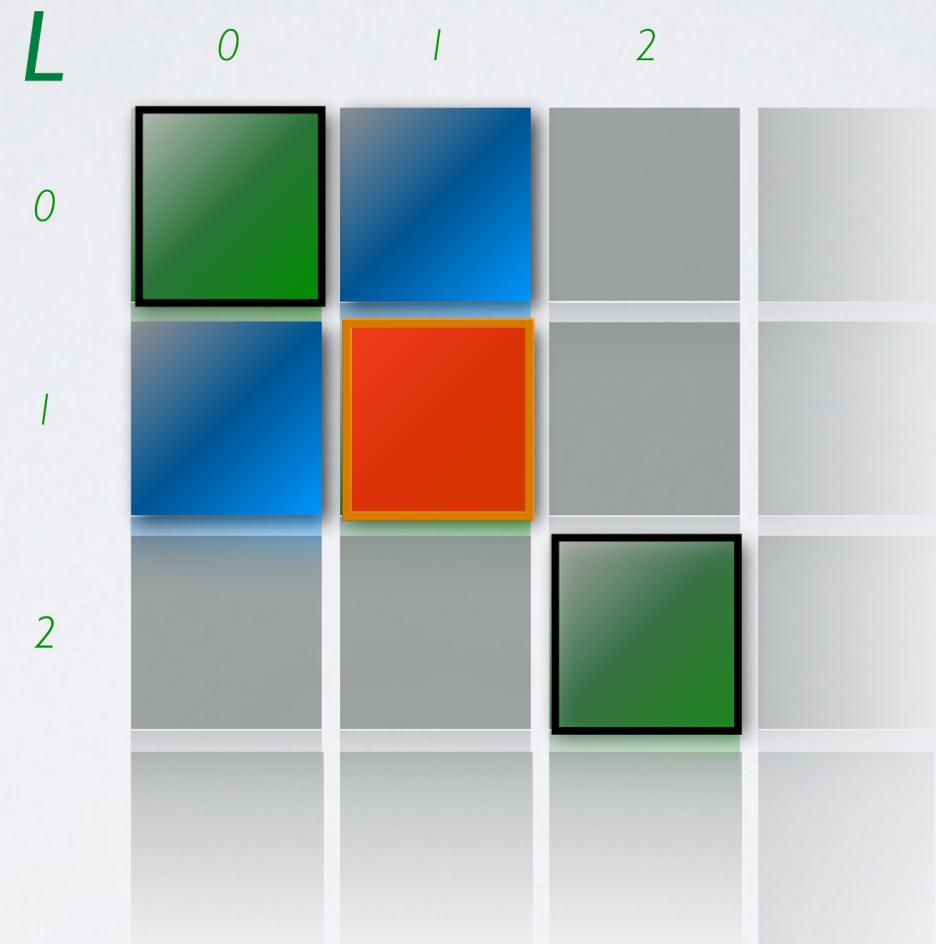
# A boundary adapted spectral approach for breakup problems

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Received 12 February 2010, in final form 11 April 2010  
Published 5 May 2010  
Online at [stacks.iop.org/JPhysB/43/101001](http://stacks.iop.org/JPhysB/43/101001)

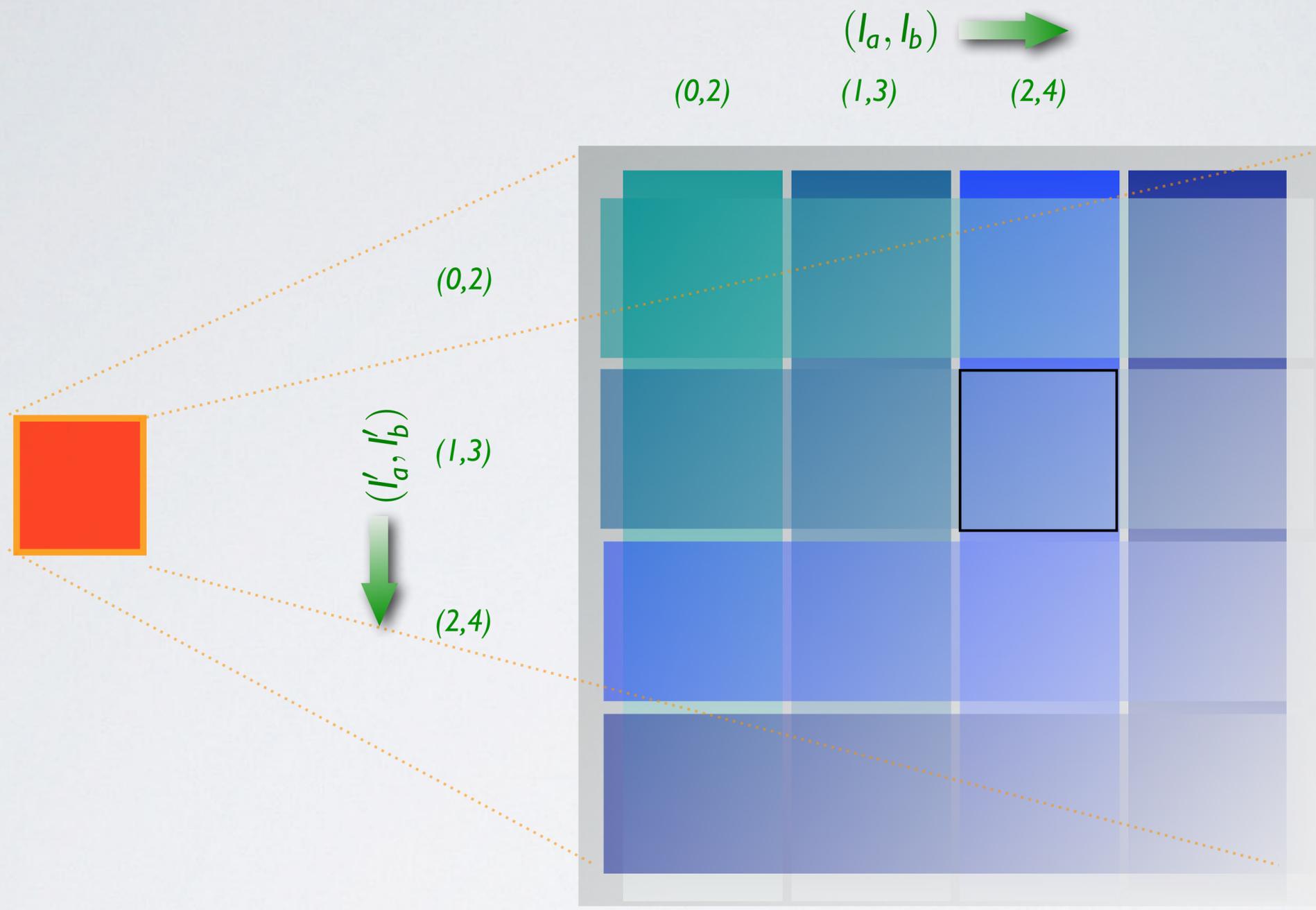
**Abstract**  
We present a spectral method to study three-body breakup processes. The method based on the Sturmian expansion and the boundary adapted Sturmian functions are solutions of one-electron Schrödinger equation with a short-range potential. Also, they are the eigenfunctions of a potential is assumed as the eigenvalue. Comparison of our analytical method demonstrates that our approach reproduces the results of the exact numerical analysis with much smaller than other previous calculations. Preliminary results of double ionization by electron impact in the  $L = 0$  approximation are discussed. Our other CI calculations, for similar basis sets, are also presented. PACS number: 31.10.+k, 31.15.Er, 31.20.+q, 31.30.+j, 31.70.+q, 31.75.+cd, 31.80.+b, 31.85.+q, 31.90.+g, 31.95.+m, 32.10.+b, 32.20.+d, 32.30.+g, 32.40.+g, 32.50.+g, 32.60.+g, 32.70.+g, 32.80.+g, 32.90.+g, 33.10.+g, 33.20.+g, 33.30.+g, 33.40.+g, 33.50.+g, 33.60.+g, 33.70.+g, 33.80.+g, 33.90.+g, 34.10.+g, 34.20.+g, 34.30.+g, 34.40.+g, 34.50.+g, 34.60.+g, 34.70.+g, 34.80.+g, 34.90.+g, 35.10.+g, 35.20.+g, 35.30.+g, 35.40.+g, 35.50.+g, 35.60.+g, 35.70.+g, 35.80.+g, 35.90.+g, 36.10.+g, 36.20.+g, 36.30.+g, 36.40.+g, 36.50.+g, 36.60.+g, 36.70.+g, 36.80.+g, 36.90.+g, 37.10.+g, 37.20.+g, 37.30.+g, 37.40.+g, 37.50.+g, 37.60.+g, 37.70.+g, 37.80.+g, 37.90.+g, 38.10.+g, 38.20.+g, 38.30.+g, 38.40.+g, 38.50.+g, 38.60.+g, 38.70.+g, 38.80.+g, 38.90.+g, 39.10.+g, 39.20.+g, 39.30.+g, 39.40.+g, 39.50.+g, 39.60.+g, 39.70.+g, 39.80.+g, 39.90.+g, 40.10.+g, 40.20.+g, 40.30.+g, 40.40.+g, 40.50.+g, 40.60.+g, 40.70.+g, 40.80.+g, 40.90.+g, 41.10.+g, 41.20.+g, 41.30.+g, 41.40.+g, 41.50.+g, 41.60.+g, 41.70.+g, 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matriz H  
bloques de L



bloque de L fijo



bloque de L fijo

$(l_a, l_b)$  →

$(0,2)$

$(1,3)$

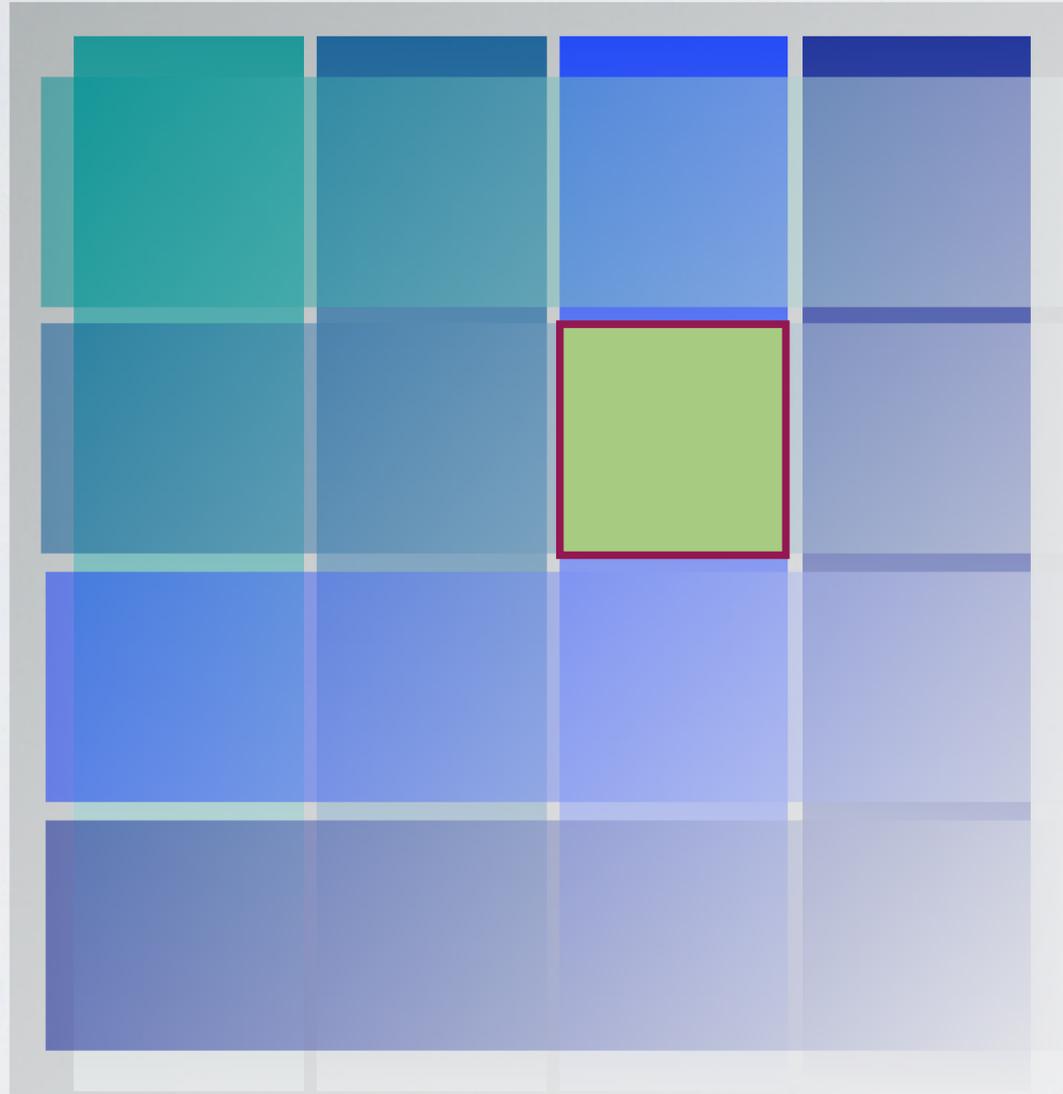
$(2,4)$

↓  $(l'_a, l'_b)$

$(0,2)$

$(1,3)$

$(2,4)$





Cada elemento de matriz H tiene



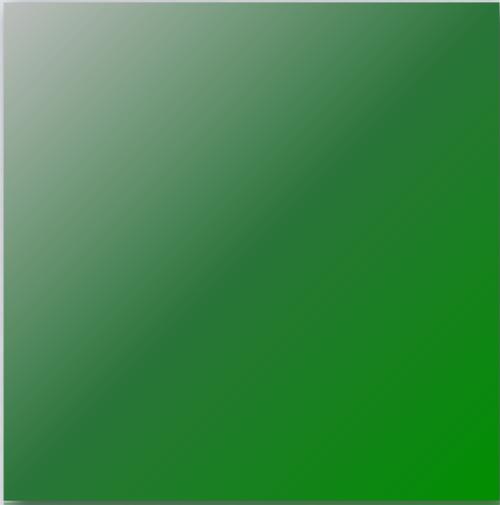
2 integrales 1D (overlaps)

prod. escalar

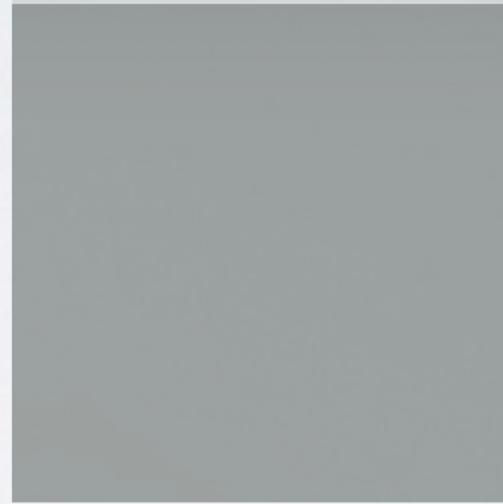
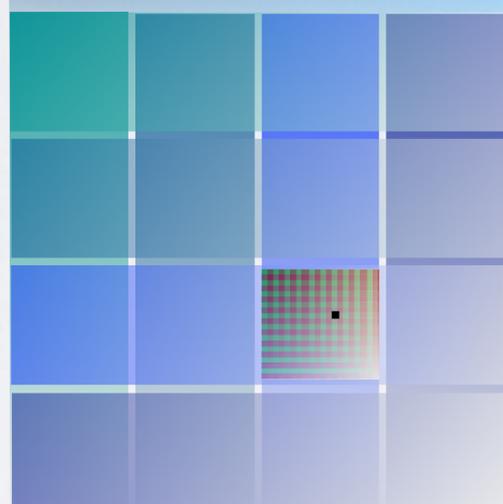
n integrales 2D (repulsión)

prefix sum  
prod. escalar

0

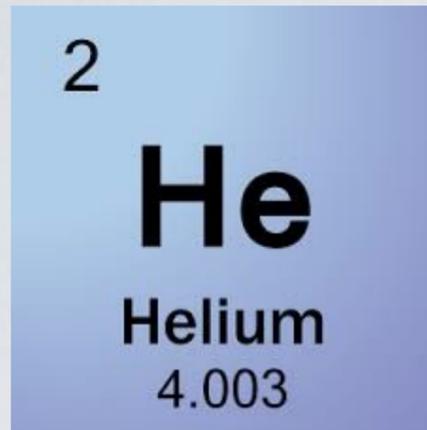


1



2



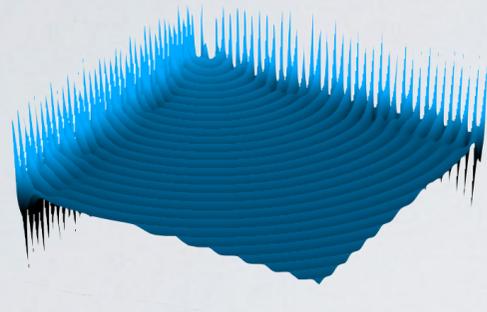


$L=0$      $\times$     13 pares  $(l_a, l_b)$      $\times$      $n=2$      $\times$      $N_a=N_b=35$

|     $\times$      $13 \times 13$      $\times$     2     $\times$      $1225 \times 1225$

253 605 625    2D

4    GB



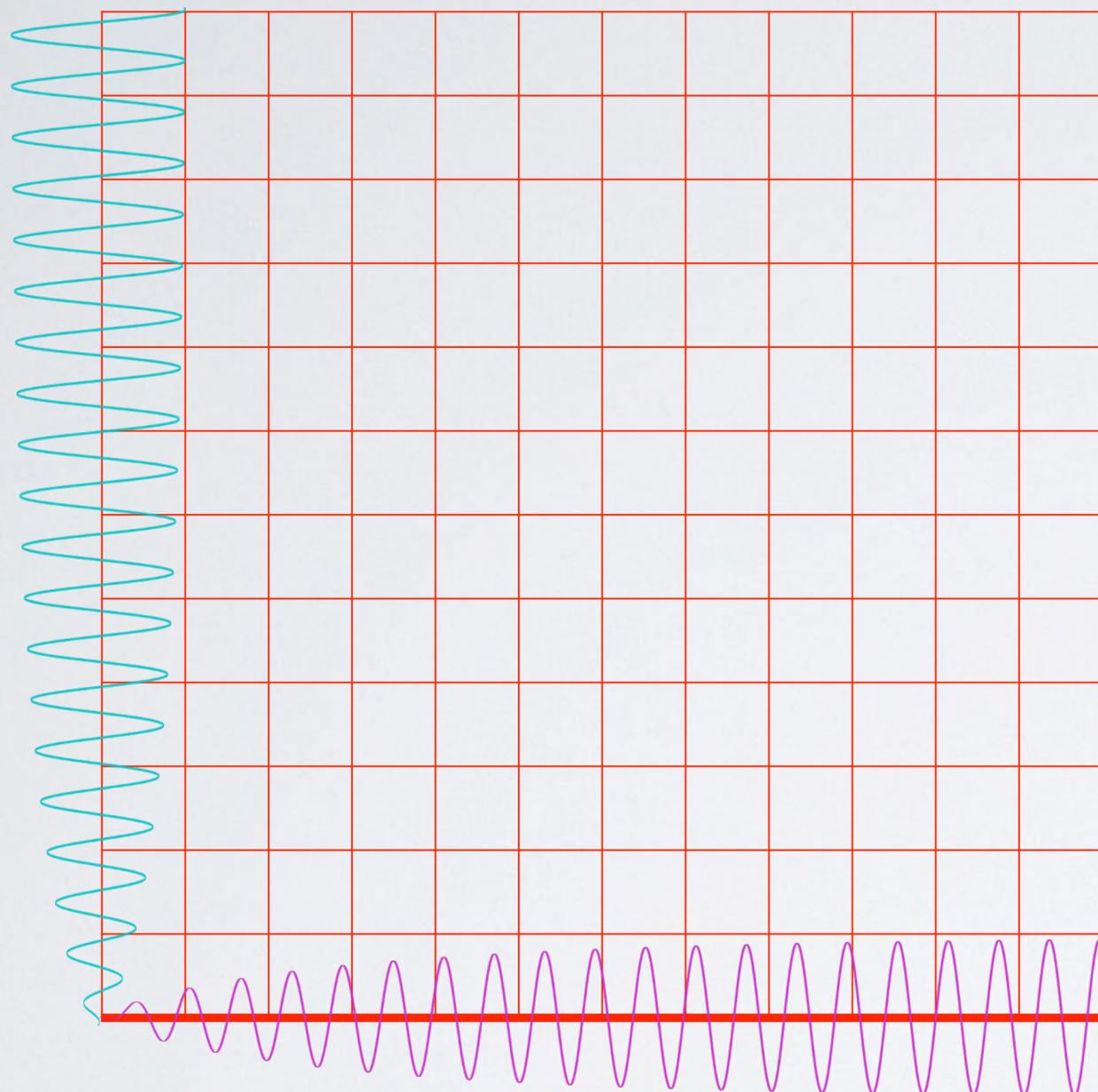
colisión (e,3e)

$L=0-4$   $\times$  5 pares  $(l_a, l_b)$   $\times$   $n=2$   $\times$   $N_a=N_b=84$

5  $\times$  5  $\times$  5  $\times$  2  $\times$  7056  $\times$  7056

12 446 784 000 2D

**98 GB**



dominio radial finito (2D) + C. C.

método espectral

cálculo de la  
matriz  $H$



problema de álgebra  
lineal numérica  
denso



cálculo de la  
matriz H



C + Fortran + CUDA

14 000 000 2D/s

25 x



## Accelerating spectral atomic and molecular collisions methods with graphics processing units



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### ARTICLE INFO

*Article history:*

Received 10 January 2014  
Received in revised form  
14 March 2014  
Accepted 26 March 2014  
Available online 3 April 2014

*Keywords:*

Ionization  
Spectral methods  
Sturmian functions  
GPU computing

### ABSTRACT

We present a computation method to accelerate the calculation of the Hamiltonian of a three-body time independent Schrödinger equation for collisions. The Hamiltonian is constructed with one dimensional (basis overlaps) and two dimensional (interparticle interaction) integrals that are mapped into a computational grid in a Graphics Processing Unit (GPU). We illustrate the method for the case of an electron impact single ionization of a two electron atom. This proposal makes use of a Generalized Sturmian Basis set for each electron, which are obtained numerically on a quadrature grid that is used to compute the integrals in the GPU. The optimal computation is more than twenty times faster in the GPU than the calculation in CPU. The method can be easily scaled to computers with several Graphics Processing Units or clusters.

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problema de álgebra  
lineal numérica  
denso



Fortran + MPI + ScaLapack

d. mitnik

north carolina (EE UU)  
metz (Francia)

cálculo de la  
matriz H

+

problema de álgebra  
lineal numérica  
denso

C + Fortran + CUDA

MPI

Fortran + MPI + ScaLapack

producción  
1er semestre 2015

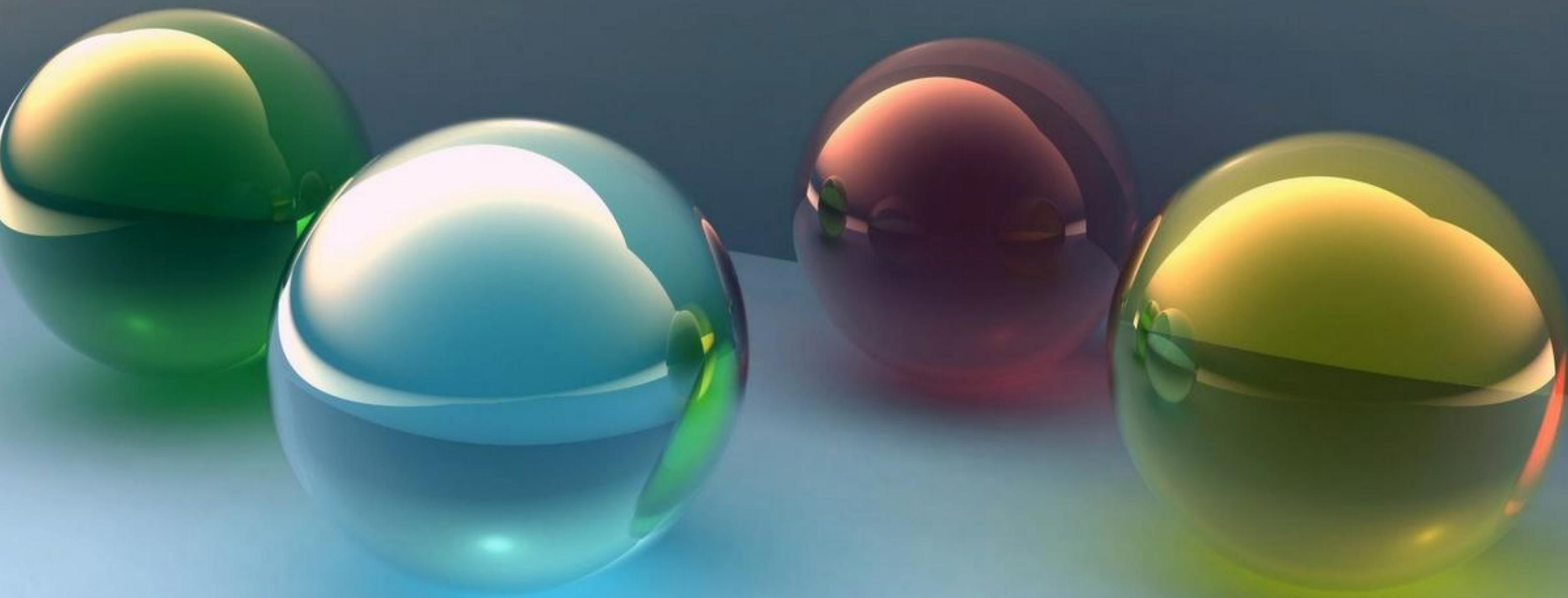






4 cuerpos

30 TB





5 cuerpos .....

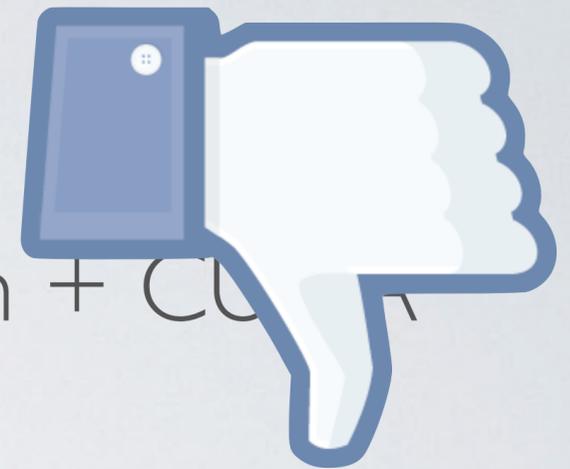
cálculo de la  
matriz H

+



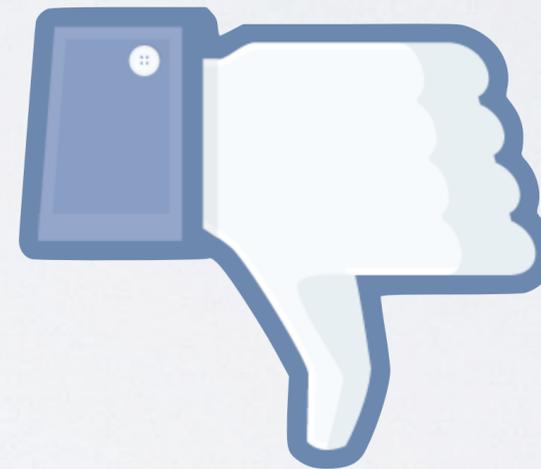
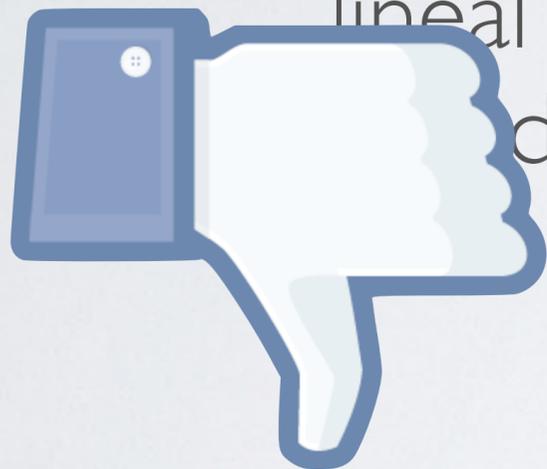
C + Fortran + CUDA

MPI



problema de álgebra  
lineal numérica  
denso

Fortran + MPI + ScaLapack



problema de álgebra  
lineal numérica  
denso

para

una matriz que no se conoce *a priori*

pero cuyos elementos se saben  
calcular

## matrices latentes

pict 2014 presentado

n wolovick

c bederián

e pilotta

l biedma

j m randazzo

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f d colavecchia

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herramientas de

HPC para  
financiación

física de colisiones y  
Agencia

reacciones químicas  
Conicet  
F. Balseiro

NVIDIA

